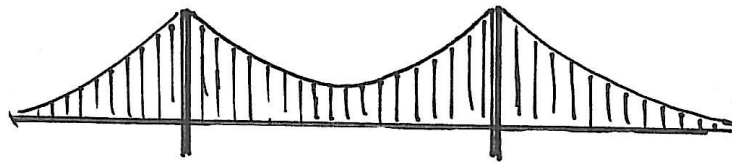


# Nuclear excitation with zeptosecond multi-MeV laser pulses

Adriana Pálffy and Hans A. Weidenmüller

Max Planck Institute for Nuclear Physics, Heidelberg, Germany

## Nuclear Fusion: From NIF to the Stars



San Francisco, August 11<sup>th</sup>, 2014



# New experimental developments ...

Experimental promise of coherent MeV photons!

**Challenge:** the theory of laser-induced nuclear reactions



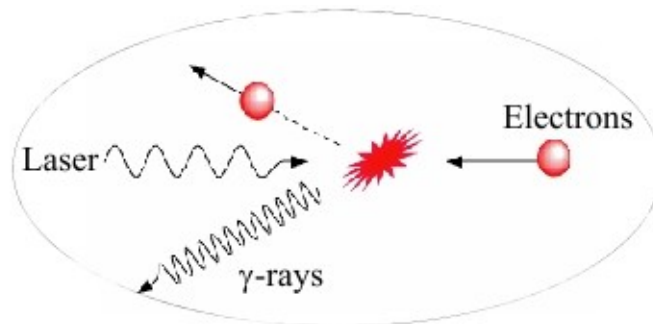
ELI beam  
(intense & optical)



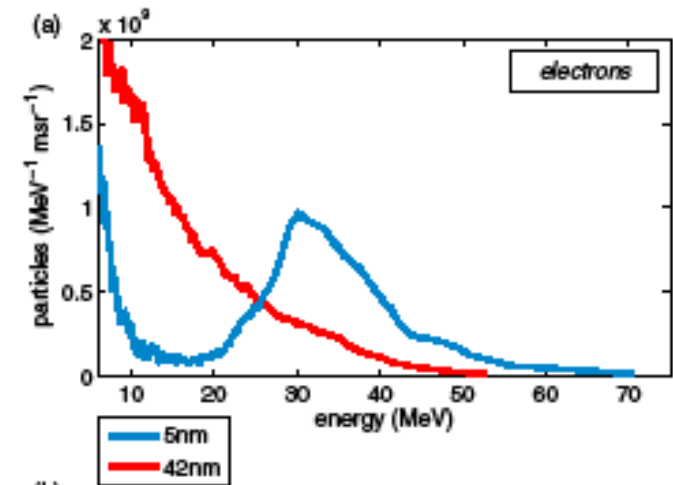
thin Carbon foil  
electron sheet



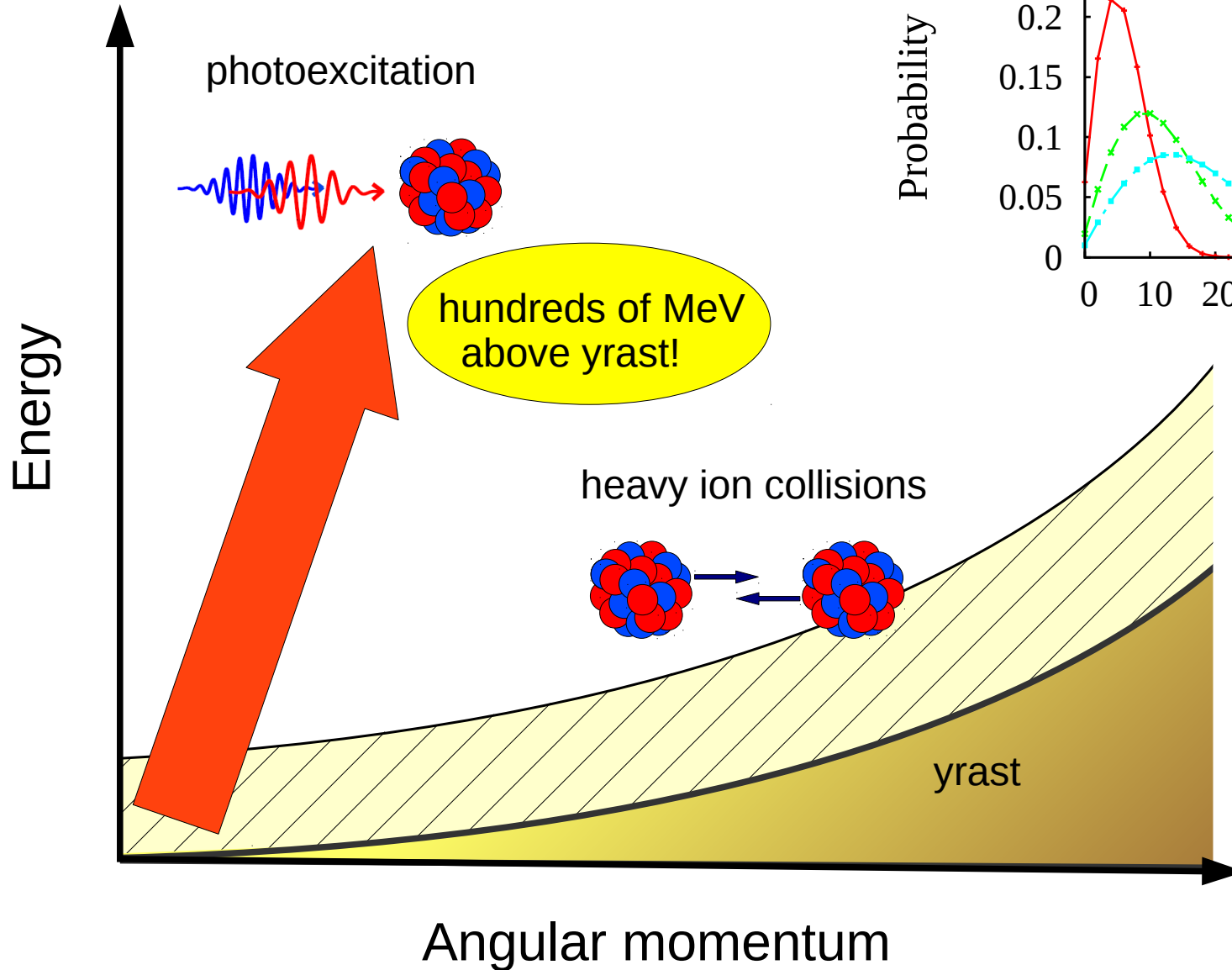
Compton backscattering  
of another laser



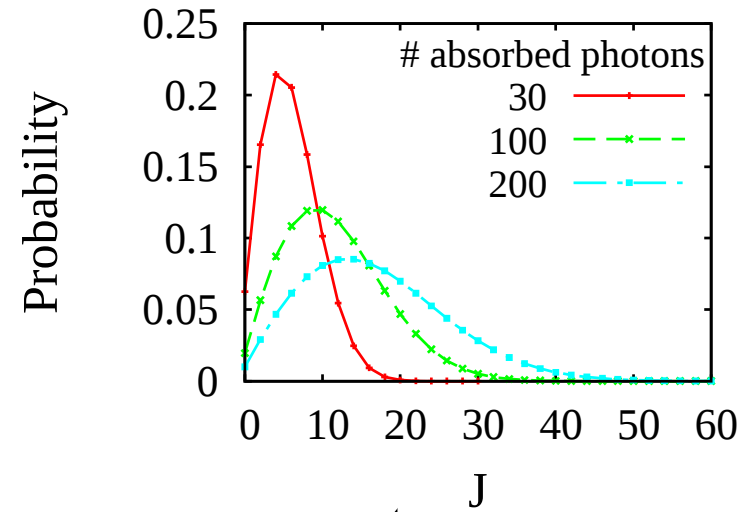
Mono-energetic & tunable  $\gamma$ -ray beam  
(unlike bremsstrahlung)



# Far from yrast



Dipole absorption  $J \sim \sqrt{N}$



# Nuclear excitation mechanism

*ELI photons*  
 $N > 10^3$   
 $E \sim 10 \text{ MeV}$   
 $T \sim 10^{-19} \text{ s}$

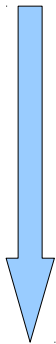


**PERTURBATIVE**

Giant Dipole Resonance once or twice

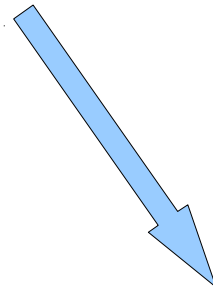
H. A. Weidenmüller, Phys. Rev. Lett. 106 (2011) 122502

B. Dietz and H. A. Weidenmüller, Phys. Lett. B 693 (2010) 316



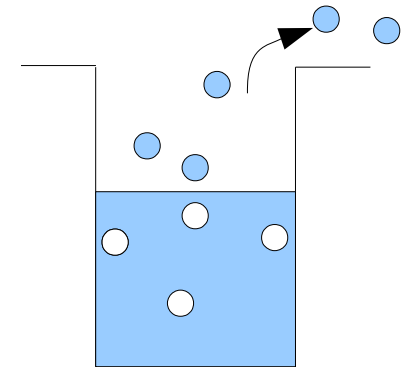
**QUASIADIABATIC**

Compound nucleus equilibrates  
about as fast as it is excited



**STRONGLY NON-ADIABATIC**

Nucleus evaporates by  
multiple nucleon emission



# Nuclear excitation mechanism

*ELI photons*  
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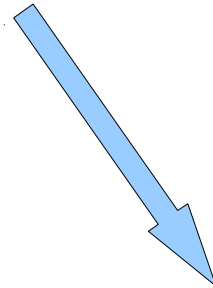
Giant Dipole Resonance once or twice

H. A. Weidenmüller, Phys. Rev. Lett. 106 (2011) 122502  
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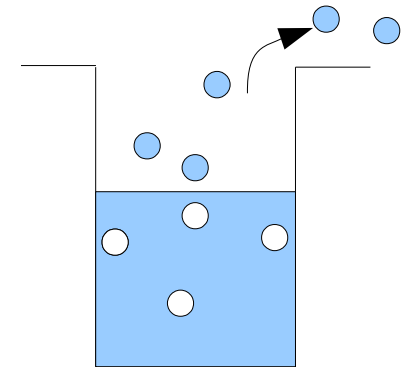
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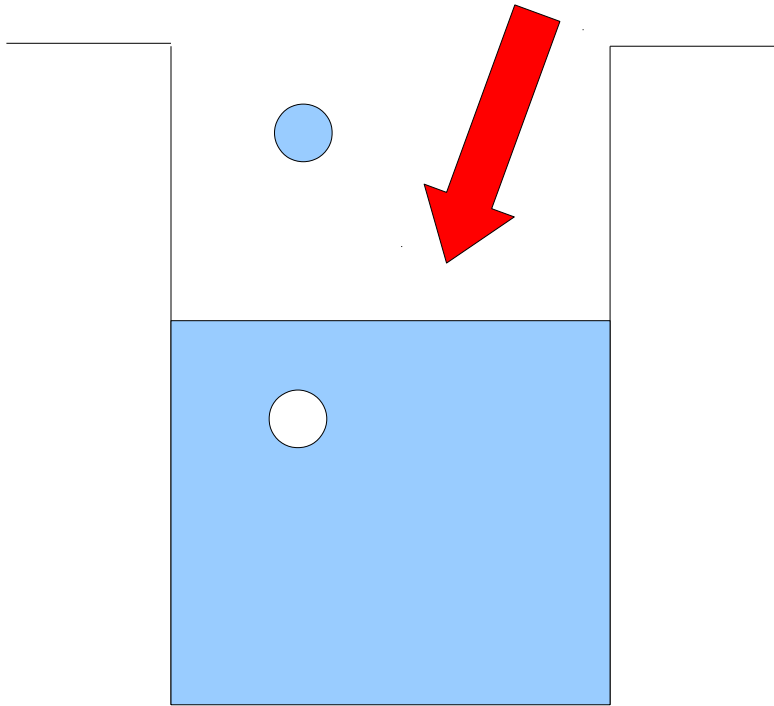


**STRONGLY NON-ADIABATIC**

Nucleus evaporates by  
multiple nucleon emission

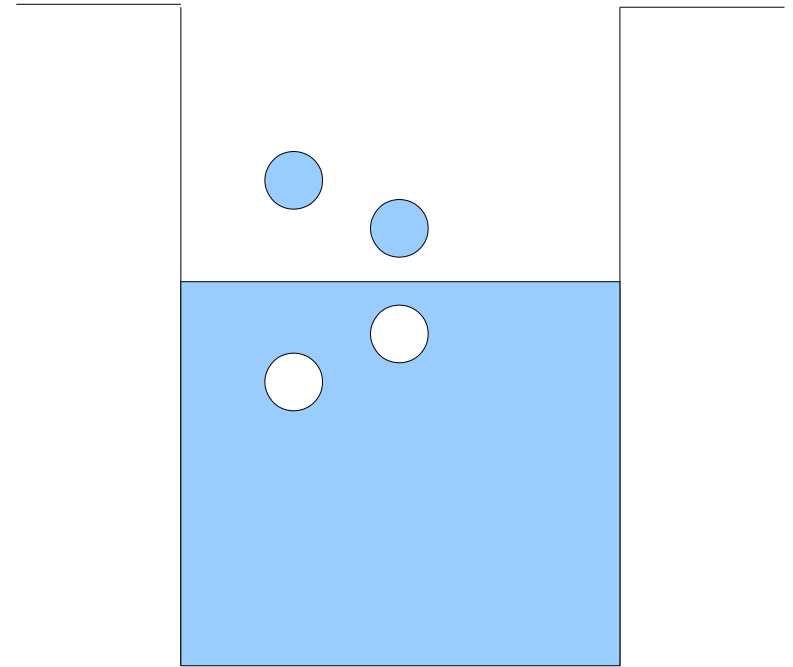


# Competing channels



PHOTOEXCITATION

creates particle-hole pairs



RESIDUAL INTERACTION

nucleon-nucleon interaction  
redistributes energy

# Competing channels

**Quasiadiabatic regime –  
after each absorption, nucleus equilibrates!!!**



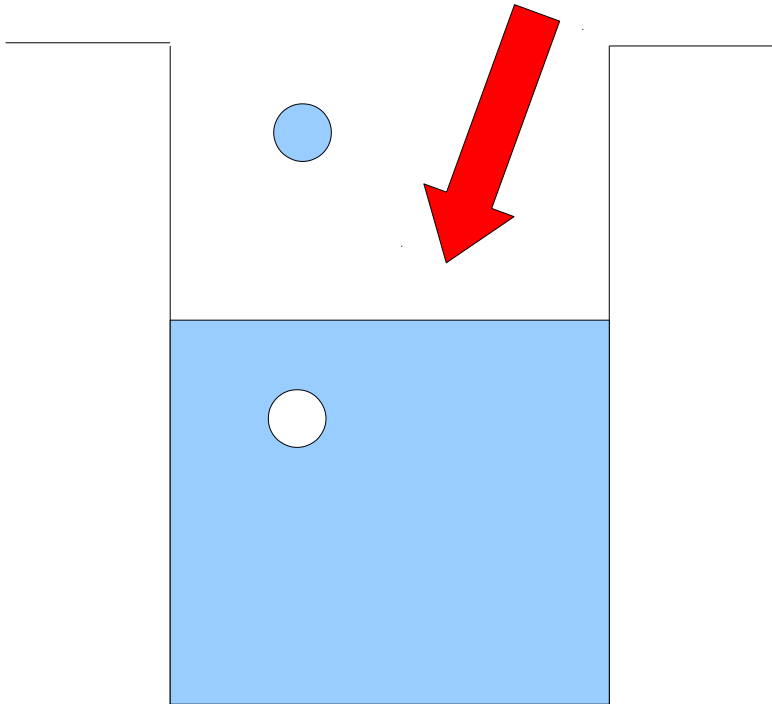
PHOTOEXCITATION

creates particle-hole pairs

RESIDUAL INTERACTION

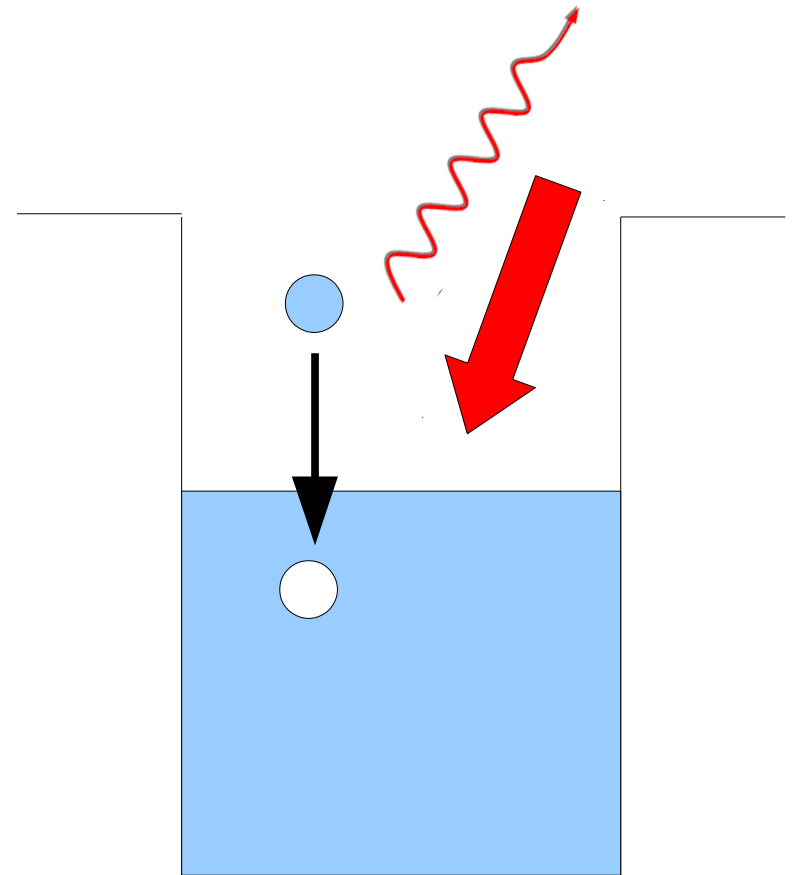
nucleon-nucleon interaction  
redistributes energy

# Competing channels



PHOTOEXCITATION

creates particle-hole pairs

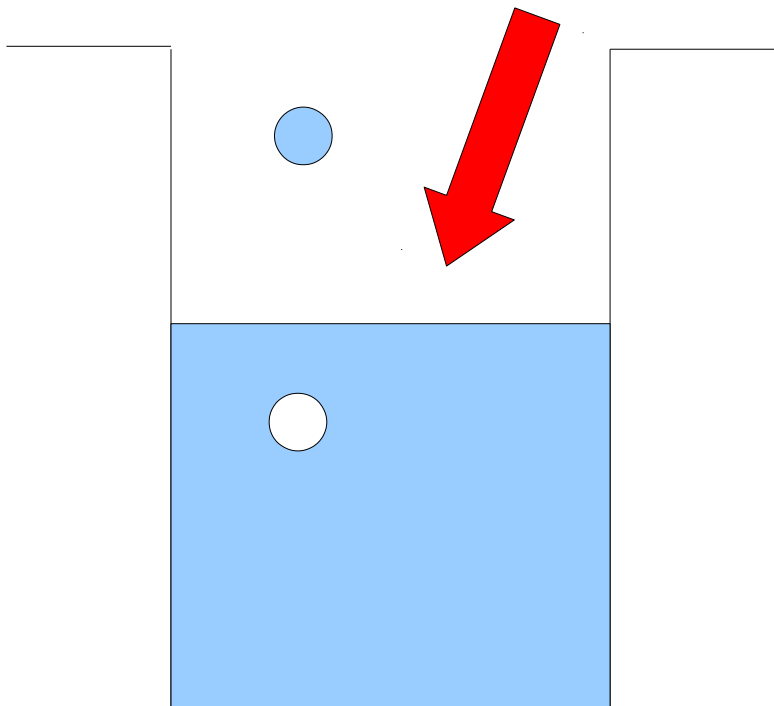


INDUCED PHOTOEMISSION

particle-hole recombination and  
photon emission

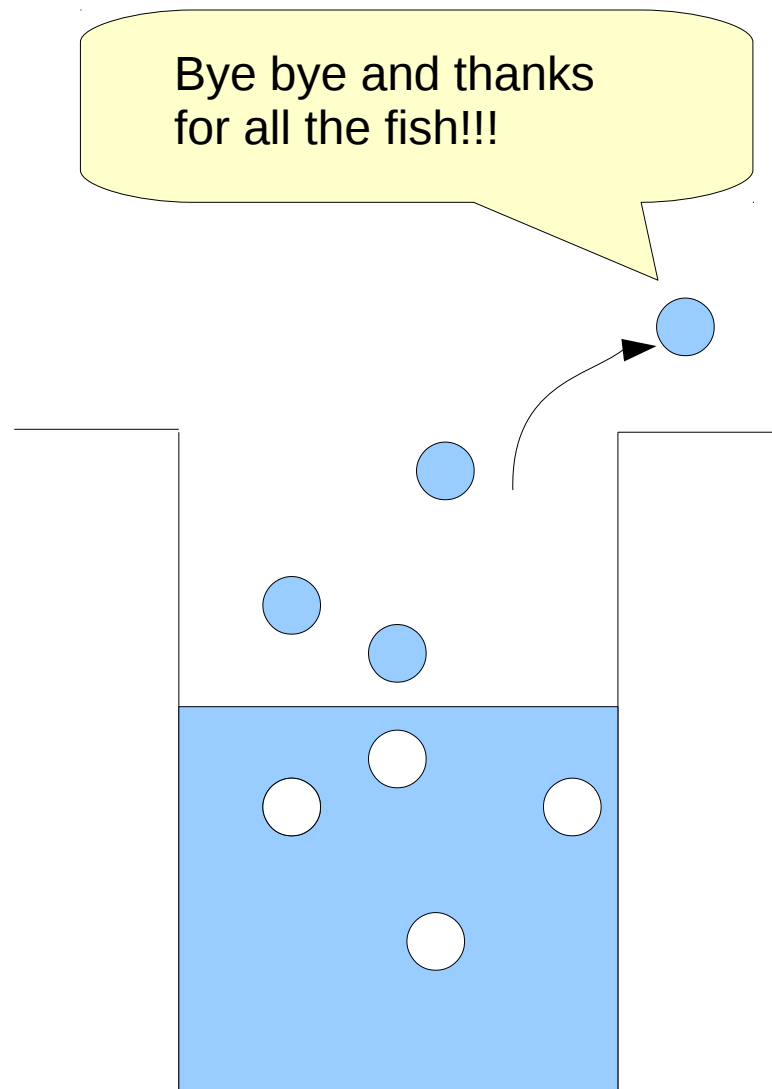


# Competing channels



PHOTOEXCITATION

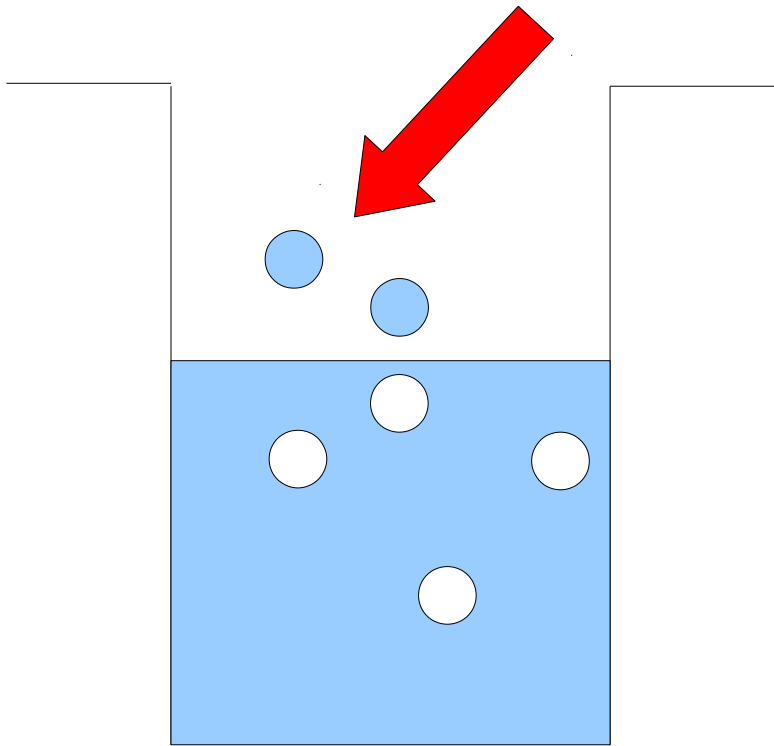
creates particle-hole pairs



NEUTRON EVAPORATION

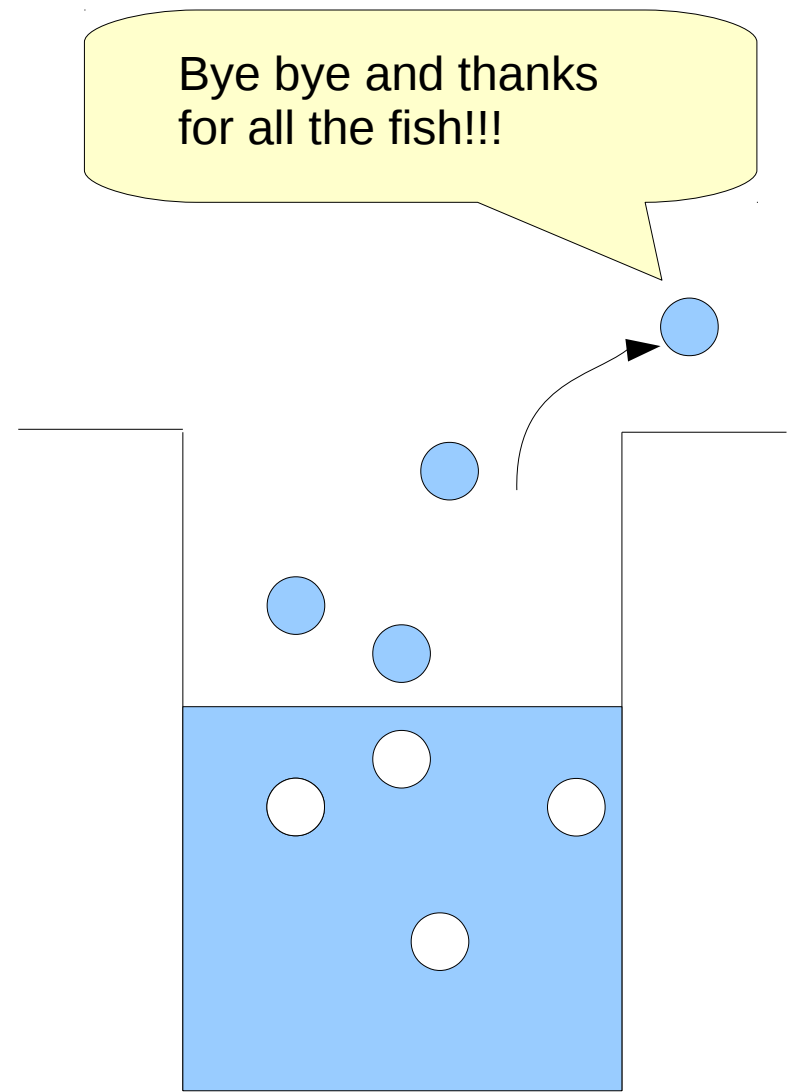
after several absorbed photons

# Competing channels



PARTICLE EMISSION

single nucleons reach the continuum



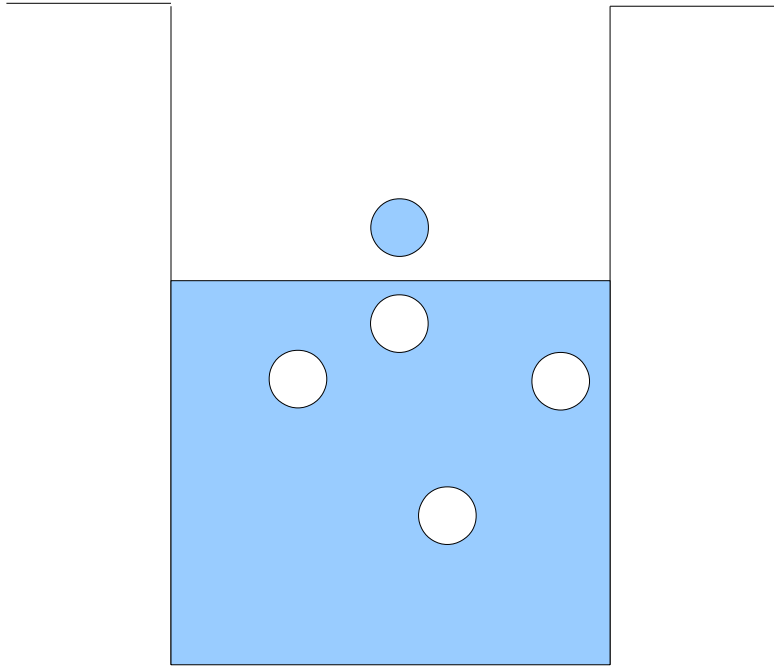
NEUTRON EVAPORATION

after several absorbed photons

# Competing channels



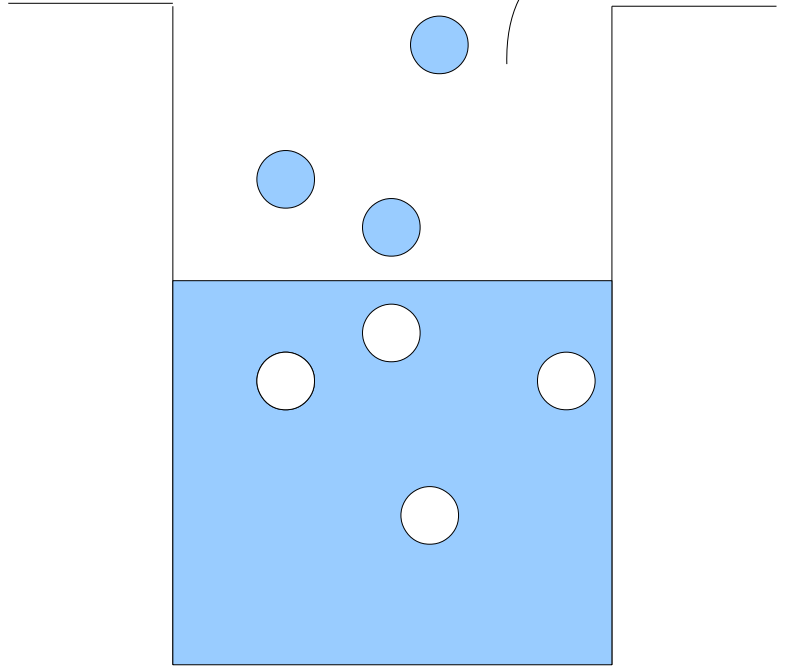
Bye bye and thanks  
for all the fish!!!



PARTICLE EMISSION

single nucleons reach the continuum

Bye bye and thanks  
for all the fish!!!



NEUTRON EVAPORATION

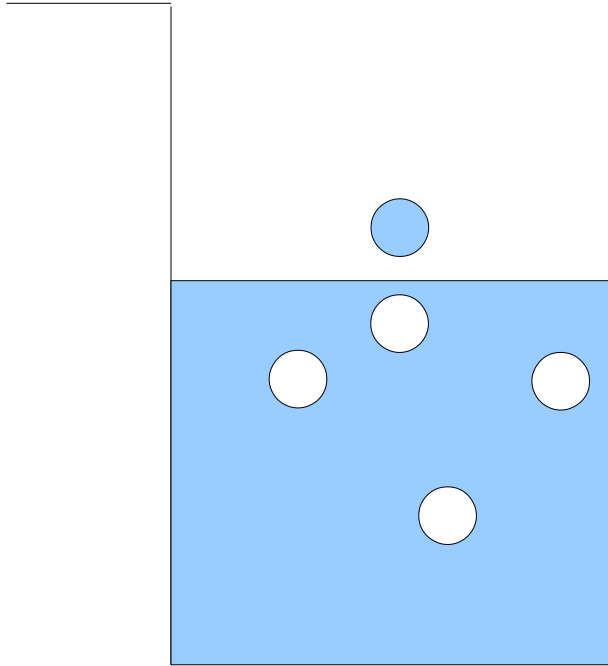
after several absorbed photons

# Competing channels

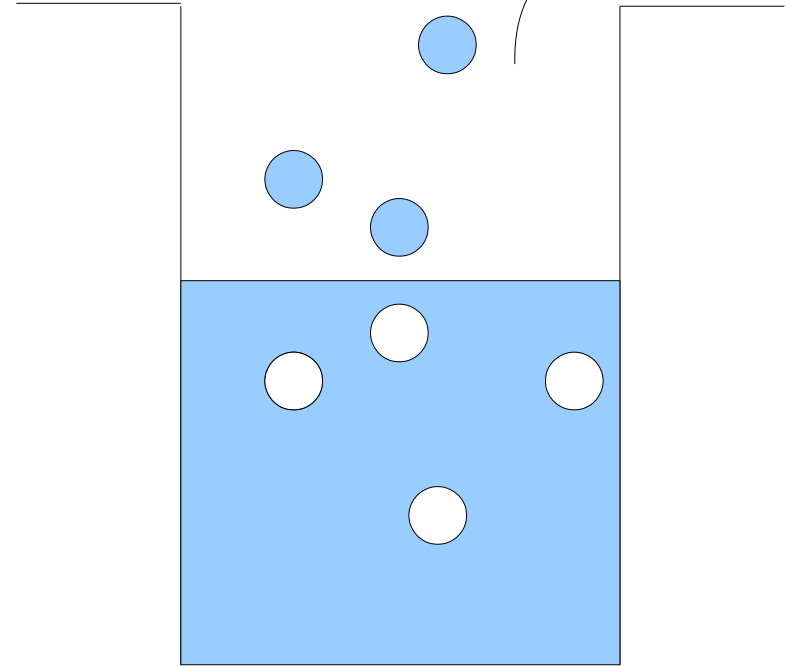


Bye bye and thanks  
for all the fish!!!

Bye bye and thanks  
for all the fish!!!



*a neutron*



PARTICLE EMISSION

single nucleons reach the continuum

NEUTRON EVAPORATION

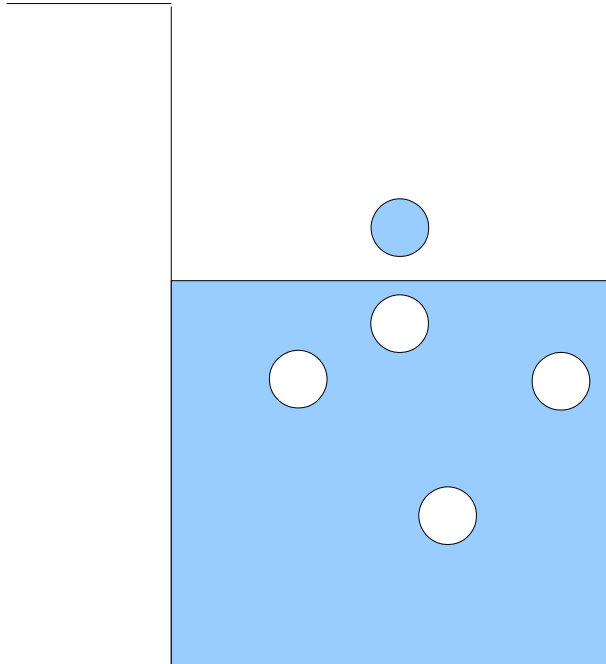
after several absorbed photons

# Competing channels

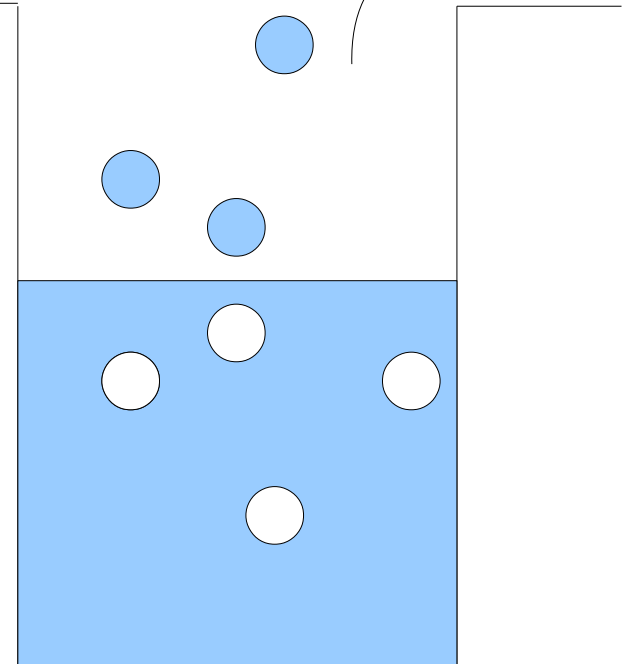


Bye bye and thanks  
for all the fish!!!

Bye bye and thanks  
for all the fish!!!



*a neutron or a proton*



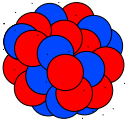
PARTICLE EMISSION

single nucleons reach the continuum

NEUTRON EVAPORATION

after several absorbed photons

# Quasiadiabatic regime



Assume complete nuclear equilibration between two photon absorptions

→ Effective absorption rate of an equilibrated compound nucleus

$$(N\Gamma)_{\text{eff}}(E) = N\Gamma_{\text{dip}} \rho_{\text{acc}}(E) / \rho_{\text{acc}}(E_g)$$

## COMPETING WITH

- Induced dipole emission

$$(N\Gamma)_{\text{ind}}(E) = (N\Gamma)_{\text{eff}}(E) \rho_A(E - E_L) / \rho_A(E)$$

- Induced nucleon emission

$$(N\Gamma)_{\text{cnt}}(E) = N\Gamma_{\text{dip}} \rho_{\text{cnt}}(E) / \rho_{\text{acc}}(E_g)$$

- Neutron evaporation

$$\Gamma_n(E) = (2\pi)^{-1} \int_{E_g(A-1)}^{E-E_n} dE' \rho_{A-1}(E') / \rho_A(E)$$

- Fission

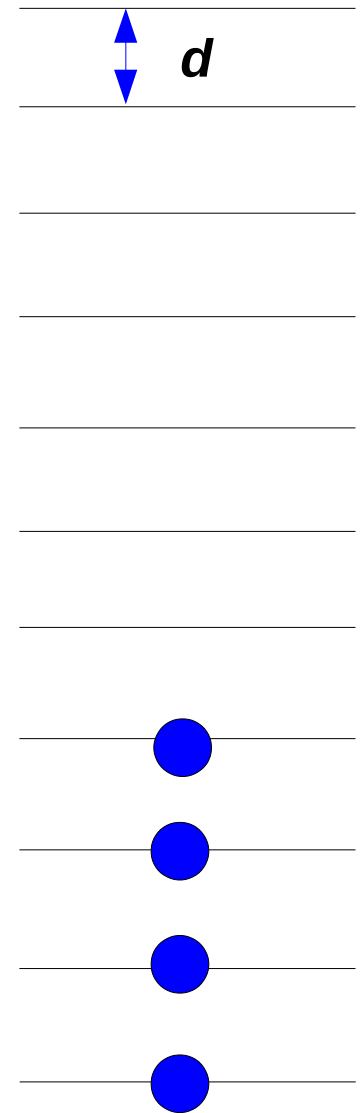
$$\Gamma_f(E) = (\hbar\omega_1 / (2\pi)) \exp\{-E_f d \ln \rho_A(E) / dE\}$$

# Level densities needed!

New theoretical formalism for high energies and high particle-hole numbers!

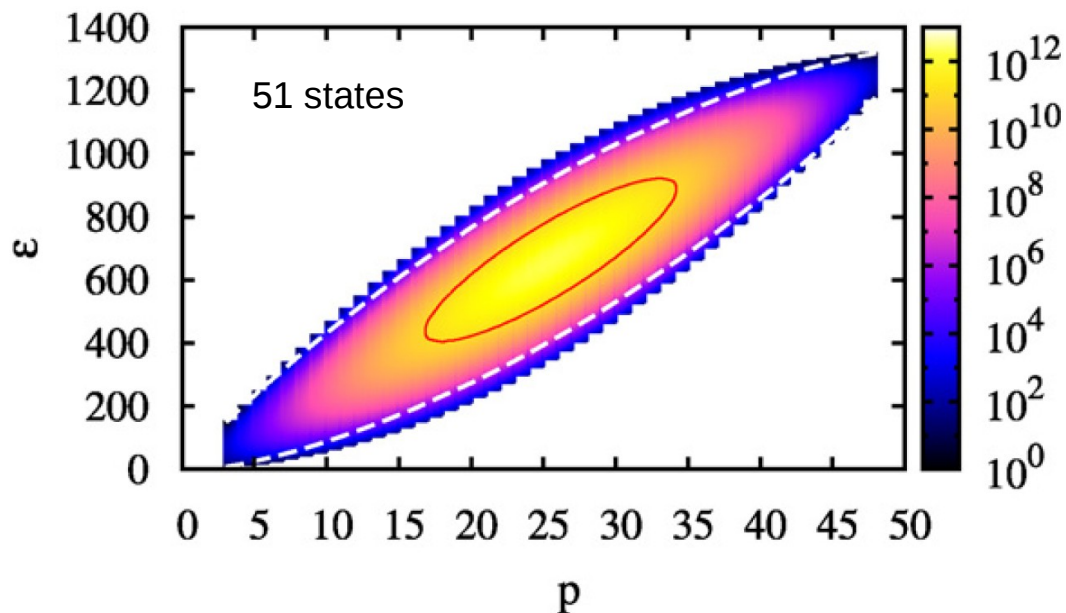
Shell model with finite number of bound states  
A spinless non-interacting fermions distributed

$$\rho_p(E, J, \pi) = \frac{1}{2} \rho_p(E) \frac{2J + 1}{2\sqrt{2\pi} \sigma_{2p}^3} \exp \left\{ - \frac{[J + (1/2)]^2}{2\sigma_{2p}^2} \right\}$$



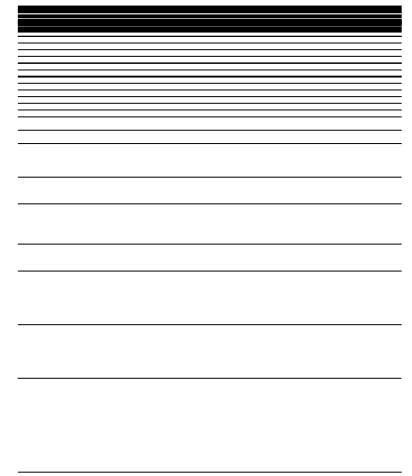
## STEP I: CONSTANT SPACING MODEL

“Corrected” Gaussian, 2<sup>nd</sup>, 4<sup>th</sup>, and 6<sup>th</sup> moments

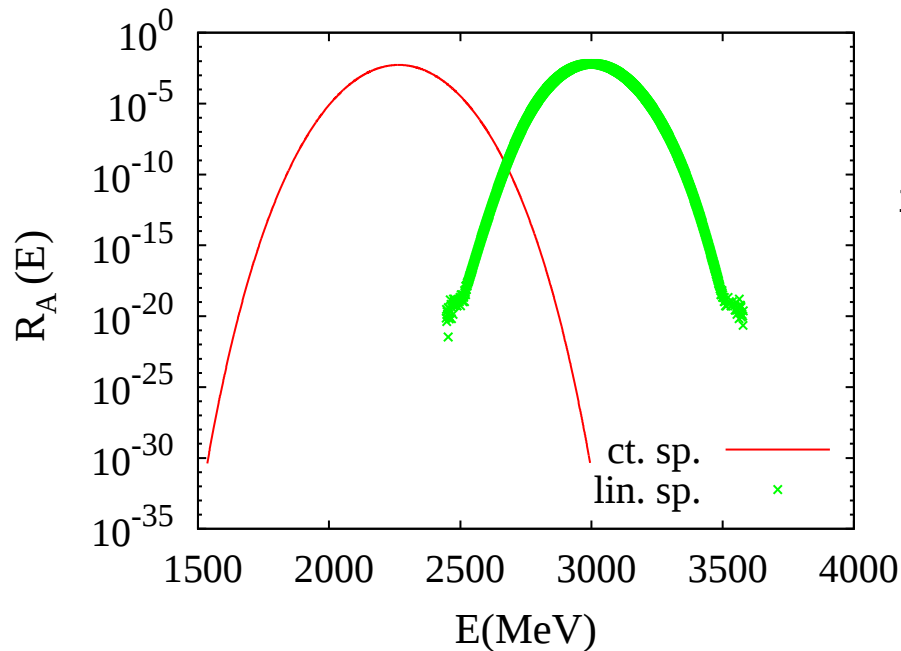


# Level densities

## STEP II: REALISTIC SPACING



More realistic case, level spacing linear or quadratic in energy



148 states  
A=100

$$\rho_1^{(1)}(\varepsilon) = \frac{2A}{F^2} \varepsilon \quad \rho_1^{(2)}(\varepsilon) = \frac{3A}{F^3} \varepsilon^2$$

Shift in energy  
Slight asymmetry  
Narrower width

AP and H. A. Weidenmüller  
Nucl. Phys. A 917, 15 (2013)

Density of accessible states - Fermi gas model



# Comparison

$$\rho_1^{(1)}(\varepsilon) = \frac{2A}{F^2} \varepsilon$$

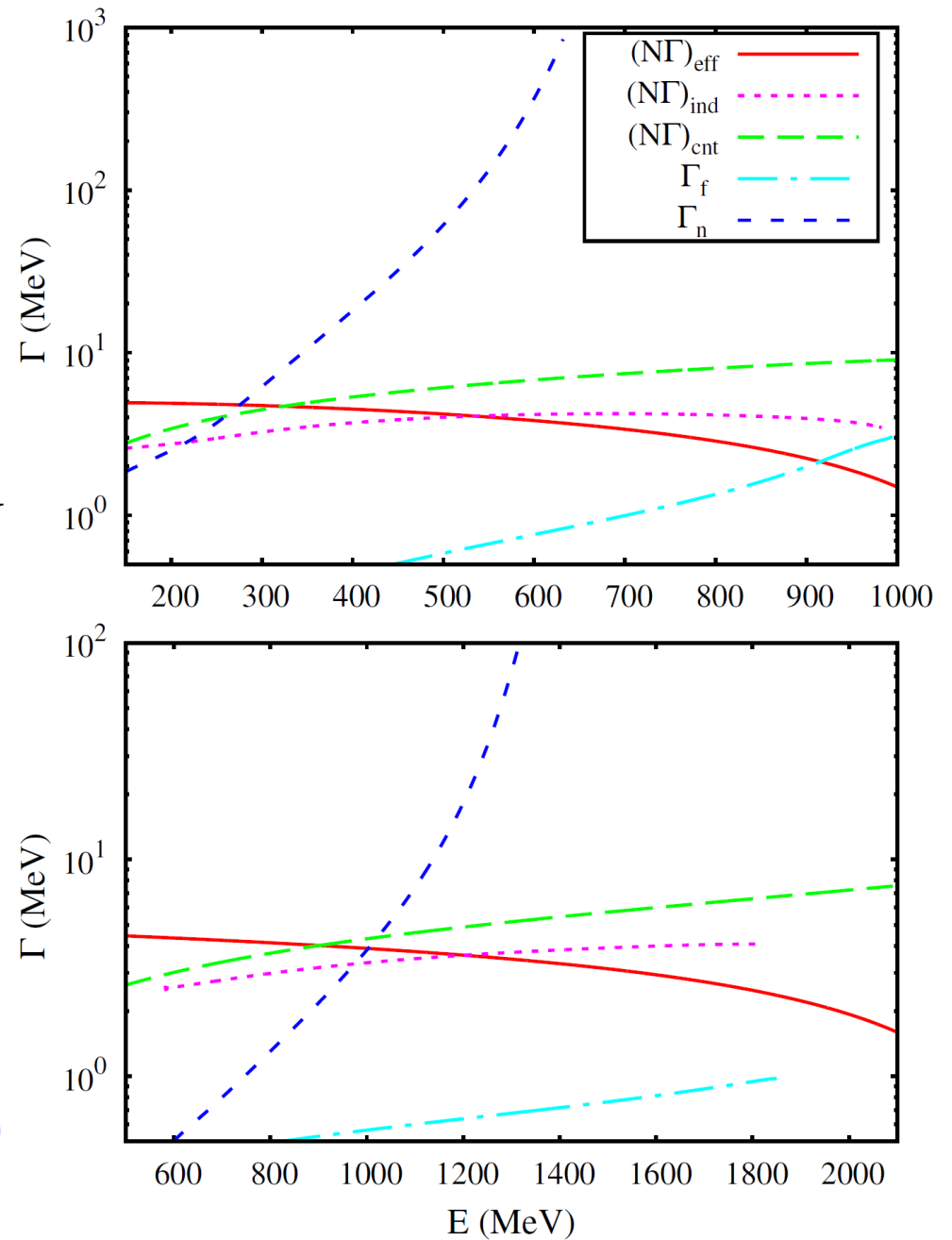
$$A = 100$$

$$N\Gamma_{\text{dip}} = 5 \text{ MeV}$$

$$\rho_1^{(2)}(\varepsilon) = \frac{3A}{F^3} \varepsilon^2$$

$$A = 200$$

AP and H. A. Weidenmüller  
Phys. Rev. Lett. 112, 192502 (2014)



!!! At maximum of level density, photon absorption and emission are equally probable!!!

# Comparison

$$\rho_1^{(1)}(\varepsilon) = \frac{2A}{F^2} \varepsilon$$

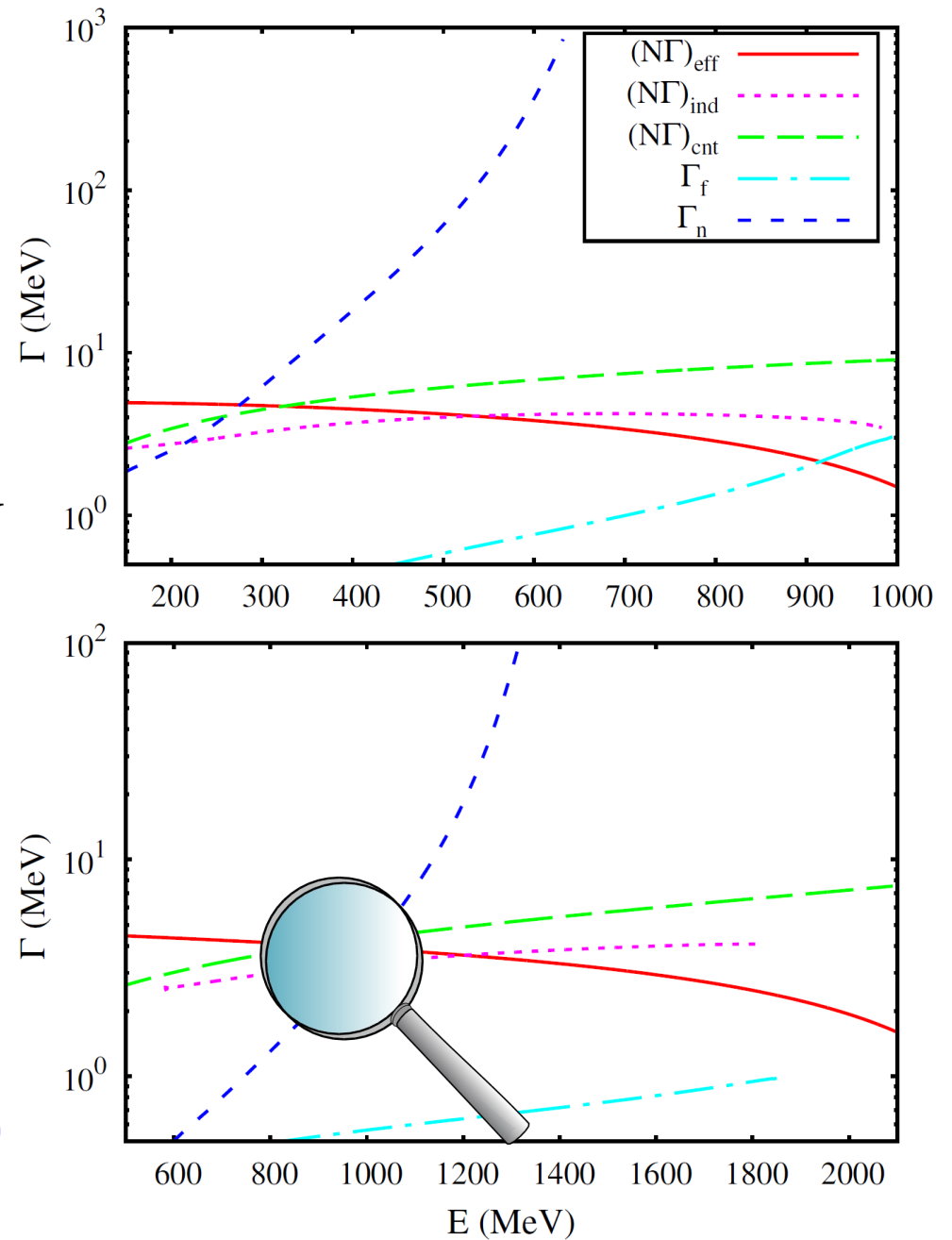
$$A = 100$$

$$N\Gamma_{\text{dip}} = 5 \text{ MeV}$$

$$\rho_1^{(2)}(\varepsilon) = \frac{3A}{F^3} \varepsilon^2$$

$$A = 200$$

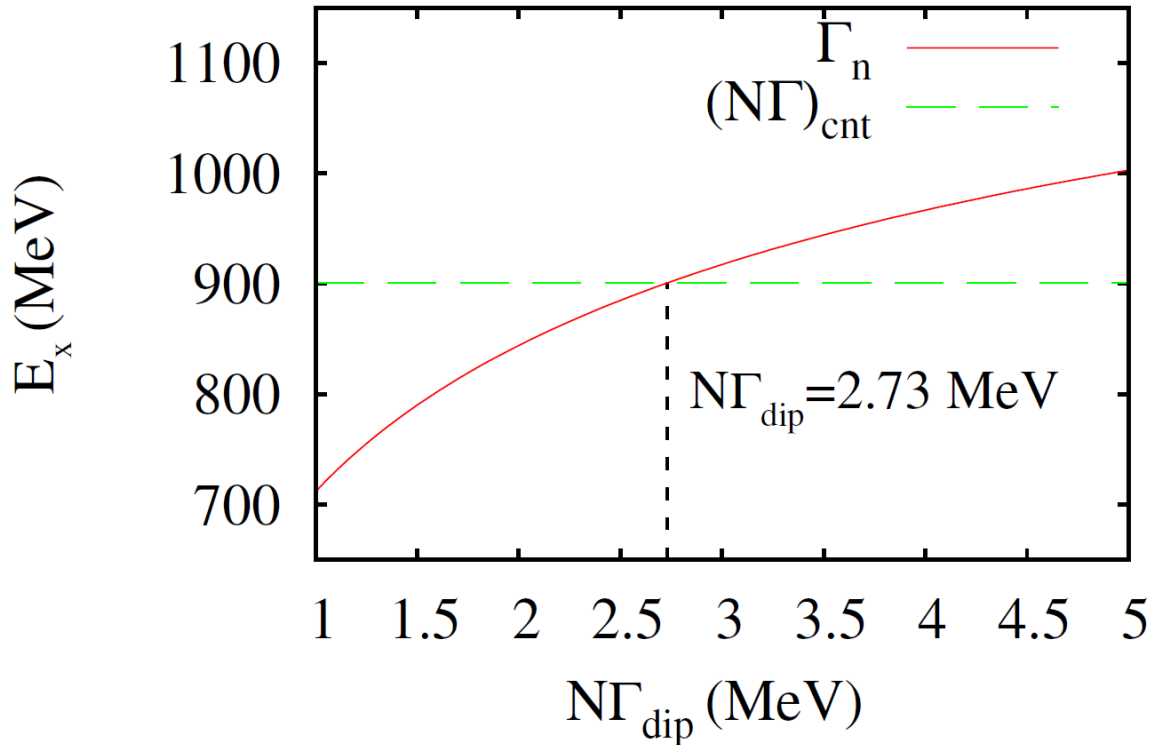
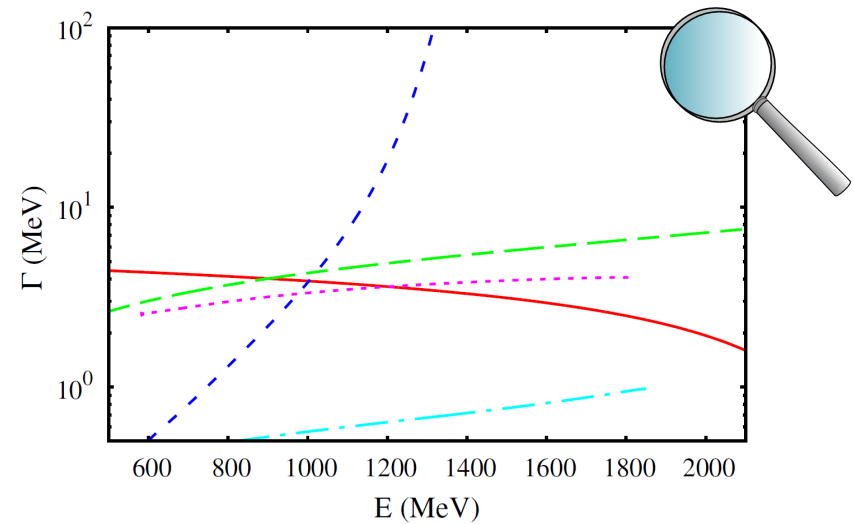
AP and H. A. Weidenmüller  
Phys. Rev. Lett. 112, 192502 (2014)



!!! At maximum of level density, photon absorption and emission are equally probable!!!

# Comparison

Emission of slow neutrons feeds states of similar energy in the daughter nuclei



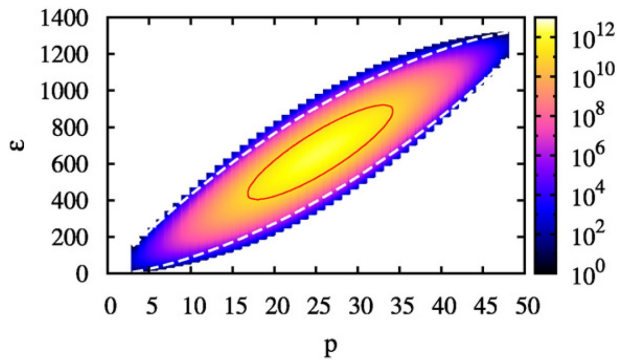
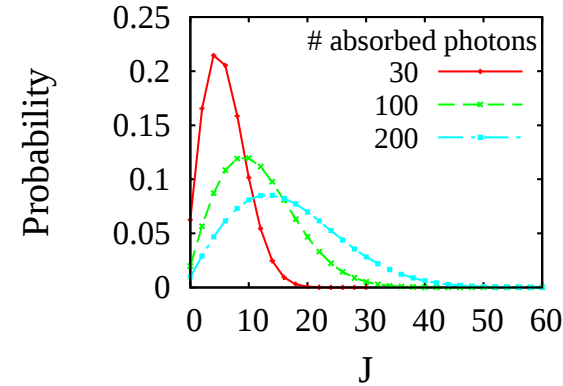
Nucleon emission – p or n - details depend on exact binding energies and absorption rates

AP and H. A. Weidenmüller  
Phys. Rev. Lett. 112, 192502 (2014)

For a 50 zs pulse, proton-rich reaction products – **FAR FROM STABILITY**

# Conclusions

→ nuclear excitation with a multi-MeV zs coherent laser pulse  
 Quasi-adiabatic regime 1 photon absorbed / nuclear relaxation time  
 leads far from yrast and far from stability!!!



→ nuclear reaction theory + newly developed method for nuclear level densities

AP and H. A. Weidenmüller  
 Phys. Lett. B 718, 1105 (2013)  
 Nucl. Phys. A 917, 15 (2013)

→ End of **reaction chain** determined by duration of laser pulse – 50 zs  
**proton-rich** nuclei due to strong neutron evaporation

AP and H. A. Weidenmüller  
 Phys. Rev. Lett. 112, 192502 (2014)

